Hans Duistermaat: the Man and his Mathematics

Johan Kolk

Mathematical Institute Utrecht University

Staff Colloquium UvA – April 2011

Thanks to Erik van den Ban Jan Brandts Wilberd van der Kallen Job Kuit V.S. Varadarajan (UCLA)

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Outline

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Biographical Data

			Biographical Data
	20-12-1942	Born in The Hague	Photographs
			Chess
►	1947-57	In the Netherlands Indies (= Indonesia)	Concentration and Persistence
	1968	Thesis at UU (= Utrecht Univ.)	Impact in Mathema
	1969-70	In Lund, Sweden, funded by Shell Prize	Mathematics Thesis
	1971	Catholic Univ. Nijmegen (= Radboud Univ.)	Proof of FIT = Fou Inversion Theorem
		, , , , , , , , , , , , , , ,	Fourier Integral Op
	1972	Full professor Nijmegen	Singularities and Eigenvalues
	1974	Professor Pure and Applied Mathematics III	Semisimple Lie Gr
	107.1	released in and and replied matternation of	Duistermaat–Heck Formula
	1982	Member KNAW (= Royal Netherlands	Classical Mechani
		Academy of Arts and Sciences)	Books
			Epiloque
	2005-09	KNAW professor. "One of the world's	Lpilogue
		most important mathematicians in differential	
		equations	
	10.00.0010		

19-03-2010 Passed away in Utrecht

Photographs



2006 Hans as we recall him from recent years.

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1977 The party following the defense of my thesis.

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Chess

Hans was an accomplished chess player. In a simultaneous match of world chess champion A. Karpov playing ten opponents in 1977, Hans was the only one to score a draw. He missed an opportunity to win.



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Everyone who met Hans, had the impression of a very powerful mind, very ascetic in temperament, extraordinarily direct and honest.

His powers of concentration and persistence were immense: in mathematics, discussions, sports,

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His powers of concentration and persistence were immense: in mathematics, discussions, sports,

For instance, when Hans wanted to go wind-surfing on the Pacific near Los Angeles. Floor, my wife took him there by car. Hans started out at 2 pm. in the Marina del Rey, an extensive harbor for yachts, and promised to be back in about 2 hours. At 5:30 Floor called me in panic, and Hans returned enthusiastically at 6:30: "it took a while to reach the ocean, but I had a wonderful time."

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Fragment of computation for a paper, to appear in Nonlinearity.

Y311 = -3 y 1 + 12 y + 22 y 1 + 312 y 4 hyperbolic Y210 = 4 1/312 - 6 y312 - 2 y311 y312 Y211 = Y321 Y212 = Y322 Y211 $Y_{321} = Y_{211} = Y_{1}^{-2} Y_{2}$ $Y_{322} = Y_{211}^{-1} Y_{212} = Y_{1}^{2} Y_{2}^{-1} Y_{1} Y_{2}^{-1} = Y_{1}^{3} Y_{2}^{-2}$ $y_1 = \frac{y_1^{-2}}{y_{321}} + \frac{y_{322}^{-1}}{y_{322}} + \frac{y_2^{-2}}{y_{321}} + \frac{y_{321}^{-4}}{y_{321}} + \frac{y_{-2}^{-2}}{y_{321}} + \frac{y_{-3}^{-2}}{y_{321}} + \frac{y_{-3}^{-2}}{y_{-3}} + \frac{y_{-3}^{-2$ $\frac{y_{301}'}{y_{301}'} = -2y_1^{-3}y_1'y_2 + \frac{y_1^{-2}y_1'}{y_1'} = -2y_1^{-3}y_2^2 + \frac{y_1^{-2}(6y_1^2 + x)}{y_1'}$ $= -2 y_{321} y_{322} y_{301} y_{392} + 6 + x y_{321} y_{322} y_{321}$ = - 2 y == + 6 + 2 y = 4 y == + 6 $\begin{array}{rcl} Y_{322} &=& 3 \, Y_1^2 \, Y_1^2 \, Y_2^{-2} \, -2 \, Y_1^3 \, Y_2^{-3} \, Y_2' \, = \, 3 \, Y_1^2 \, Y_2^{-1} \, -2 \, Y_1^3 \, Y_2^{-3} \, \left(6 \, Y_1^2 \, + \, \infty \right) \\ &=& 3 \, Y_{321}^{-4} \, Y_{322}^{-2} \, 3 \, X_{321}^2 \, Y_{322}^{-12} \, -12 \, Y_{321}^{-10} \, Y_{322}^{-5} \, Y_{321}^{-9} \, Y_{322}^{-2} \, -2 \, \alpha \, Y_{321}^{-6} \, Y_{322}^{-6} \, y_{321}^{-6} \, Y_{322}^{-6} \end{array}$ $= 3 y_{201}^{-1} - 12 y_{301}^{-1} y_{322}^{-1} - 29 y_{201}^{3} y_{300}^{3}$

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u(0)=0 u'(0)=1 u'(0)=0 u(0)=0 $(3) u(3) = 2 + 3^{3} (-\frac{1}{02} + (2-p)w(3))$ What is the mindual value of peak + psesix (-++ + (peak - p) w/peip)) over all w(x) = 2 wn =" with radius of convergence > p = $pe^{ix} + p^3 e^{3ix} \left(-\frac{1}{4\pi} + p \sum_{n=0}^{\infty} p^{ny} e^{ih + ihx} - p \sum_{n=0}^{\infty} p^n e^{inx}\right)$ $= p e^{i\alpha} - \frac{p^3}{n^2} e^{3i\alpha} + \sum_{n \neq q} p^{n+q} e^{i(n+q)\alpha} - \sum_{n \neq q} p^{n+3} w_n e^{i\xi}$ $\Rightarrow u_{0} = 0 \quad \inf_{w,p} h_{p} \left[1 - p^{2} e^{2im} \left(-p^{-2} + \left(p e^{im} - p \right) w \left(p e^{im} \right) \right) \right]$ $p = bpe^{ih}$ $|1 + e^{ih}(-b^2 e^{2ib} + (e^{ih} + be^{ih})p^2 u(pe^{ih}))|$ +) + 10 - b) 210 16 3 400 4.= 0 un 11+ (Pe") (-1+ (Pe"-1) pw/pe")) 1+==2(-1+(3-1)w(3)) tor w(z) = -020 2(8-4 1160.0 1/(0) = 0 and then azelight persentation 0- close to P+ Geluri 011100 $(\rho^2 + \rho_+^2) E d d = + \rho_+^3 - \rho_-^3$ Mad 2. × 11937) 671-57 1 12K1 1g(2) >1 (P==P=)[(P=+P+)= -(P=+P=P==)] pomber in sources P= D= (2-P=) $(P_{-}P_{+})(P_{+}+P_{+})$, $(P_{-}+P_{+})f(P_{+}+P_{+})P_{+} - (P_{-}^{2}+P_{-}P_{+}+P_{+}^{2})$ (2) g(2) (b(2) (5) P2P= (=-P+) (a)=zg(8) (c)=0 pla \$=0: +-9=1/6 1-2001572.15.2001 Manperson --9700 WAY

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Brilliant research mathematician in geometric analysis.

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Brilliant research mathematician in geometric analysis.

Active in different parts of math

- Ordinary and partial differential equations
- Discrete integrable systems (difference equations, complete sets of conservation laws, exact solutions)
- Classical mechanics
- Analysis on (semisimple) Lie groups (= smooth manifolds that are (special) groups. Example: certain groups of matrices)
- Symplectic differential geometry (= differential geometry describing the phase spaces of classical mechanical systems and generalizations)
- Algebraic geometry

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- 23 PhD students and furthermore post-docs
- International contacts, in particular, USA, France, Indonesia
- Inspiring teacher

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Author citations in Math Reviews Citation Database 1071 Duistermaat

For comparison 2079 A. Schrijver 1594 N.Yu. Reshetikhin 7618 J.-P. Serre (Fields medalist) 667 J.-C. Yoccoz (Fields medalist)

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Video

Rothschild Visiting Professor Discrete Integrable Systems programme Isaac Newton Institute for Mathematical Sciences Cambridge, U.K. 18-05-2009

http://www.sms.cam.ac.uk/media/603139

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Thesis

Energy and entropy as real morphisms for addition and order. Foundations of thermodynamics. Originated in seminar conducted by Günther K. Braun, professor of applied mathematics, and a chemist.

Braun died tragically one year before defense of thesis. Thesis advisor: Hans Freudenthal, geometer.

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Thesis

Energy and entropy as real morphisms for addition and order. Foundations of thermodynamics. Originated in seminar conducted by Günther K. Braun, professor of applied mathematics, and a chemist.

Braun died tragically one year before defense of thesis. Thesis advisor: Hans Freudenthal, geometer.

Thesis lead to dissent between physicists and mathematicians at UU.

Hans dropped the topic.

Nevertheless an important influence in his career.

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Thesis

Studied works of Sophus Lie on contact and symplectic transformations (modern terminology).

Lars Hörmander, analyst and Fields medalist, used Hans' knowledge in formulation of calculus of FIO's = Fourier integral operators.

This class of operators contains partial differential operators as well as classical integral operators as special cases, while composition of these operators is often well-defined.

One goal is to obtain inverses of differential operators. Useful in the theory of high-frequency waves and propagation of singularities, like shock waves.

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Fourier transform \mathcal{F} acting on Schwartz function ϕ on \mathbf{R}^n (= smooth and rapidly decaying at infinity)

$$\mathcal{F}\phi(\xi) = \int_{\mathbf{R}^n} e^{-i\langle x,\xi\rangle} \phi(x) \, dx = \int_{\mathbf{R}^n} e_{-i\,\xi} \, \phi \, dx.$$

 $\mathcal{F}\phi$ is Schwartz too.

FIT = Fourier Inversion Theorem FIO

$$\phi(\mathbf{x}) = \frac{1}{(2\pi)^n} \int_{\mathbf{R}^n} e^{i\langle \mathbf{x}, \xi \rangle} \mathcal{F} \phi(\xi) \, d\xi.$$

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Fourier transform \mathcal{F} acting on Schwartz function ϕ on \mathbf{R}^n (= smooth and rapidly decaying at infinity)

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 $\mathcal{F}\phi$ is Schwartz too. FIT = Fourier Inversion Theorem **FIO**

$$\phi(x) = \frac{1}{(2\pi)^n} \int_{\mathbf{R}^n} e^{i\langle x,\xi\rangle} \mathcal{F}\phi(\xi) \,d\xi.$$

Sufficient to prove for x = 0:

$$(2\pi)^n \phi(\mathbf{0}) = \int_{\mathbf{R}^n} \mathcal{F}\phi(\xi) \, d\xi$$

Indeed, integration by parts leads to local and global intertwining identities, for $\partial_j := \partial/\partial x_j$, **•**

$$\mathcal{F} \circ \partial_j = i \, \xi_j \circ \mathcal{F}, \qquad \mathcal{F} \circ T_a^* = e_{i \, a} \circ \mathcal{F}, \quad T_a^* \phi(x) = \phi(x + a).$$

For general case of FIT, replace ϕ by $T_x^*\phi$.

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Desired: $(2\pi)^n \phi(0) = \int_{\mathbf{R}^n} \mathcal{F}\phi(\xi) d\xi$. Use Dirac's delta function δ and formally interchange order of integration to rewrite this as

$$(2\pi)^n \int_{\mathbf{R}^n} \delta(x) \, \phi(x) \, dx = \int_{\mathbf{R}^n} \int_{\mathbf{R}^n} e^{-i\langle x,\xi \rangle} \, d\xi \, \phi(x) \, dx.$$

In other words, FIT equivalent to

$$(2\pi)^n \delta = \int_{\mathbf{R}^n} \boldsymbol{e}_{-i\,\xi} \, \mathbf{1} \, d\xi = \mathcal{F}\mathbf{1}.$$

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In other words, FIT equivalent to

$$(2\pi)^n \delta = \int_{\mathbf{R}^n} \boldsymbol{e}_{-i\xi} \, \mathbf{1} \, d\xi = \mathcal{F} \mathbf{1}.$$

Tempered distributions are linear functionals acting on Schwartz functions continuously with respect to a suitable topology. Examples: $\delta(\phi) = \phi(0)$, $\psi(\phi) = \int_{\mathbf{R}^n} \psi(x) \phi(x) dx$.

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$$(2\pi)^n \delta = \int_{\mathbf{R}^n} \boldsymbol{e}_{-i\xi} \, \mathbf{1} \, d\xi = \mathcal{F} \mathbf{1}.$$

Tempered distributions are linear functionals acting on Schwartz functions continuously with respect to a suitable topology. Examples: $\delta(\phi) = \phi(0)$, $\psi(\phi) = \int_{\mathbf{R}^n} \psi(x) \phi(x) dx$. For tempered distribution *T*, define $(\mathcal{F}T)(\phi) = T(\mathcal{F}\phi)$. Justification: for Schwartz function *T* we have

$$(\mathcal{F}T)(\phi) = \int_{\mathbf{R}^n} (\mathcal{F}T)(\xi) \, \phi(\xi) \, d\xi = \int_{\mathbf{R}^n} T(\xi) \, \mathcal{F}\phi(\xi) \, d\xi = T(\mathcal{F}\phi).$$

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For years discussions on and off about most transparent proof of FIT.

Hans came up with following one-line argument:

$$\mathbf{0} = \mathcal{F}\mathbf{0} = \mathcal{F}(\partial_j \mathbf{1}) = i\,\xi_j\,\mathcal{F}\mathbf{1} \qquad (\mathbf{1} \leq j \leq n).$$

Indeed, follows from local intertwining identity.



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Indeed, follows from local intertwining identity. Argument implies that $T := \mathcal{F}1$ is supported at $\{0\}$. Therefore action of T on smooth functions is well-defined.

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Indeed, follows from local intertwining identity. Argument implies that $T := \mathcal{F}1$ is supported at $\{0\}$. Therefore action of T on smooth functions is well-defined. Next, first-order Taylor expansion gives $\phi = \phi(0) \mathbf{1} + \sum_{j} \xi_{j} \psi_{j}$. Now $\xi_{j} T = 0$ leads to existence of $c \in \mathbf{C}$ with

$$\mathcal{F}\mathbf{1}(\phi) = T(\phi) = \phi(\mathbf{0}) T\mathbf{1} + \sum_{j=1}^{n} \xi_j T(\psi_j) = \mathbf{c} \,\delta(\phi).$$

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Testing identity against Gaussian $e^{-\|\cdot\|^2/2}$ gives $c = (2\pi)^n$. Thus FIT \bullet follows.

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Fourier Integral Operators

For $P(x, \partial)$, linear partial differential operator with variable coefficients, preceding results imply

$$P(x,\partial)\phi(x) = \frac{1}{(2\pi)^n} \int_{\mathbf{R}^n} e^{i\langle x,\xi\rangle} P(x,i\xi) \mathcal{F}\phi(\xi) d\xi$$

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$$P(x,\partial)\phi(x) = \frac{1}{(2\pi)^n} \int_{\mathbf{R}^n} e^{i\langle x,\xi\rangle} P(x,i\xi) \mathcal{F}\phi(\xi) d\xi$$

= $\frac{1}{(2\pi)^n} \int_{\mathbf{R}^n \times \mathbf{R}^n} e^{i\langle x-y,\xi\rangle} P(x,i\xi) \phi(y) d(y,\xi)$

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Proof of FIT = Fourier nversion Theorem

Fourier Integral Operators

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Fourier Integral Operators

For $P(x, \partial)$, linear partial differential operator with variable coefficients, preceding results imply

$$P(x,\partial)\phi(x) = \frac{1}{(2\pi)^n} \int_{\mathbf{R}^n} e^{i\langle x,\xi\rangle} P(x,i\xi) \mathcal{F}\phi(\xi) d\xi$$

= $\frac{1}{(2\pi)^n} \int_{\mathbf{R}^n \times \mathbf{R}^n} e^{i\langle x-y,\xi\rangle} P(x,i\xi) \phi(y) d(y,\xi)$
= $\frac{1}{(2\pi)^n} \int_{\mathbf{R}^n \times \mathbf{R}^n} e^{i\chi(x,y,\xi)} P(x,i\xi) \phi(y) d(y,\xi).$

Relax conditions on phase function χ and on *P* to get FIO's.

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Singularities and Eigenvalues

Influential articles arising from the calculus of FIO's

- ▶ (jointly with Hörmander) FIO's II (1972)
- Oscillatory integrals, Lagrange immersions and unfolding of singularities (1974)
 Agenda for study of singularities of smooth functions and their applications to distribution theory.
 In some sense complementary to FIO's and parallel to work of Vladimir Arnol'd, Dr.h.c. UU
- (jointly with Victor Guillemin) The spectrum of positive elliptic operators and periodic bicharacteristics (1975) Application of FIO's to the asymptotics of spectra (= distribution of eigenvalues) of elliptic operators. Relation between spectra and geometry. His most cited paper <

These firmly established Hans' mathematical reputation.

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Singularities and Eigenvalues

Inventiones math. 29, 39-79 (1975) © by Springer-Verlag 1975

The Spectrum of Positive Elliptic Operators and Periodic Bicharacteristics

J.J. Duistermaat (Utrecht) and V.W. Guillemin (Cambridge, Mass.)

Introduction

Let X be a compact boundaryless C^{*} manifold and let P be a positive elliptic self-adjoint pseudodifierential operator of order m > 0 on X. For technical reasons we will assume that P operates on half-densities rather than functions. (We will denote the half-density bundle over X by Ω_{\pm}) We will also assume that P is a classical pseudodifferential operator in the sense that on every coordinate patch its total symbol $\sigma_p(x, \xi)$ admits an asymptotic expansion

$$\sigma_p(x,\xi) \sim \sum_{j=0}^{\infty} p_{m-j}(x,\xi)$$

with $p_{m-j}(x, \xi)$ homogeneous of degree m-j. We recall that the principal symbol p of P is equal to p_m on local coordinates, and the subprincipal symbol is equal to $p_{m-1} - \frac{1}{2i} \sum \frac{\partial^2 p}{\partial x_j \partial \xi_j}$.

Let $\lambda_1, \lambda_2, ...$ be the eigenvalues of *P*. It was remarked by Chazarain in [6] and by ourselves in [11] that the sum $\sum e^{-i \sqrt{T_{kf}}}$ is well-defined as a generalized function of *t* and that if *T* is in its singular support then the Hamiltonian vector field

$$H_q = \frac{\partial q}{\partial \xi_j} \frac{\partial}{\partial x_j} - \frac{\partial q}{\partial x_j} \frac{\partial}{\partial \xi_j}, \quad q = \sqrt[m]{p}$$

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Transition to Lie theory was remarkable occurrence. It demanded complete change of mathematical perspective: from classical and highly general aspects of PDE's to very specific questions on semisimple Lie groups.

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Transition to Lie theory was remarkable occurrence. It demanded complete change of mathematical perspective: from classical and highly general aspects of PDE's to very specific questions on semisimple Lie groups.

Improve error estimates in \bigcirc , in special case of compact locally symmetric space *X*, to prove conjectures of Israel Gel'fand. Such an *X* is described in terms of a semisimple Lie group *G*.

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DKV1 (V = V.S. Varadarajan) uses Selberg trace formula, deep generalization of Poisson's summation formula on \mathbf{R} :

$$\sum_{\lambda \in \mathbf{Z}} \mathbf{e}_{i\,\lambda} = 2\pi \sum_{n \in \mathbf{Z}} \delta_{2\pi\,n}, \qquad \sum_{\lambda \in \mathbf{Z}} \mathcal{F}\phi(\lambda) = 2\pi \sum_{n \in \mathbf{Z}} \phi(2\pi n).$$

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Here $X = \mathbf{R}/2\pi \mathbf{Z}$ and LHS (spectral) in terms of eigenvalues of the Laplace operator $-\partial^2$ on *X*. RHS (geometric) in terms of vol(*X*) and lengths of periodic geodesics of *X*.

Simultaneous control of behavior of elementary spherical functions $x \mapsto \phi_{\lambda}(x)$ on *G* as both λ and *x* to infinity independently of each other, by method of stationary phase.

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Nowhere pointwise convergence.

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Semisimple Lie Groups

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It was in the second paper (DKV2) that a true synthesis was achieved, of ideas that Hans was truly in love with, such as oscillatory integrals, caustics, Lagrange immersions, with the ideas in semisimple Lie theory, such as orbital integrals and their asymptotics.

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It was in the second paper (DKV2) that a true synthesis was achieved, of ideas that Hans was truly in love with, such as oscillatory integrals, caustics, Lagrange immersions, with the ideas in semisimple Lie theory, such as orbital integrals and their asymptotics.

Highly original suggestion to Van den Ban. Take the integrals on real flag manifolds that represent elementary spherical functions ϕ_{λ} , and move them into the complex flag manifolds,

by moving the real cycle into the complex domain.

A generalization of classical technique in complex function theory of deforming the contour,

but now in a geometrically more complex multi-dimensional context.

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Duistermaat-Heckman Formula

Gert Heckman wrote thesis under G. van Dijk and Hans on questions concerning irreducible representations of compact Lie groups.

Gert and Hans proved these results could be understood in context of symplectic geometry.

This led to celebrated Duistermaat–Heckman formula It expresses that Fourier transform of canonical (Liouville) measure on a symplectic manifold is given exactly by stationary phase approximation.

Second most quoted article.

It led to many important generalizations and is frequently used in new applications.

Some years ago famous mathematician J.-M. Bismut declared that he spent most of his life to understand the meaning of this formula.

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Classical mechanics

- Periodic solutions of a Hamiltonian system near an equilibrium point
- Monodromy in integrable Hamiltonian systems
- Nonholonomically constrained systems. Examples: disk or ball with center of mass not at its geometric center, rolling without slipping on horizontal plane under influence of constant vertical gravitational force.

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Books

Hans (co-)authored a total of 11 books (with Richard Cushman, Wiktor Eckhaus, Kolk, Jędrzej Śniatycki). 3 books were published last year.

- Lie Groups, 2000. Most cited
- Fourier Integral Operators, 1973, 1996, 2011.
 Second most cited

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Multidimensional Real Analysis II Integration

J. J. Duistermaat, J. A. C. Kolk

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Hans Duistermaat

Books

Reviews

- Fourier Integral Operators
 Duistermaat's notes have certain features that make them
 unique
- The Heat Kernel Lefschetz Fixed Point Formula for the Spin-c Dirac Operator
 Overall this is a carefully-written, highly readable book on a very beautiful subject
- Lie Groups A brilliant book on Lie groups
- Multidimensional Real Analysis
 The mother of all multivariable calculus books
- Distributions: Theory and Applications This is a fantastic book.

... I am ecstatic to have the book by Duistermaat and Kolk in my hands, as it is a gem

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Graphics Epilogue

Graphics

Intense interest in correct graphics. Hand-colored picture of SL(2, R).



Hans Duistermaat

Graphics

Graphics

Dot for dot on a matrix printer. Image of character of **SU**(3).



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Graphics

Pencil of real sine-Gordon curves, with eight real base points.



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For much more on Hans' scientific work, consult

- Nieuw Arch. Wiskd. (5) 11 (2010), 235–243
- Geometric Aspects of Analysis and Mechanics. A Conference in Honor of the 65th Birthday of Hans Duistermaat, 2011 to be published by Birkhäuser, Boston
- Notices Amer. Math. Soc., Obituaries, to appear
- http://www.staff.science.uu.nl/~kolk0101/ duisinnem html

A truly remarkable Renaissance mind

(V.S. Varadarajan)

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Epilogue

WONDER, Wiskunde onderzoekschool Nederland (= The Dutch Research School in Mathematics) has decided to establish the

Hans Duistermaat Chair for Visiting Professors



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